
Haakon Andresen

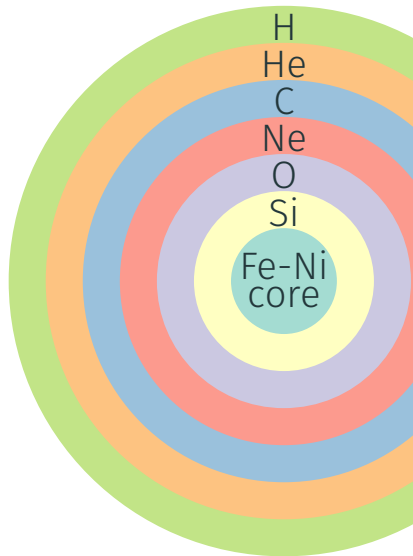
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Stellar Evolution

Shell burning

- Stars of initial masses above $\sim 8M_{\odot}$
- Shell burning
- Degenerate iron core



Cores must be massive enough to avoid becoming degenerate

- Radiation dominated: $P = \frac{1}{3}aT^4$
- Ideal gas: $P = k\rho T$
- Non-relativistic degenerate electrons: $P = K(\rho/\mu_e)^{\frac{5}{3}}$
- Degenerate electrons: $P = K(\rho/\mu_e)^{\frac{4}{3}}$

Stars in hydrostatic equilibrium:

$$P_c = \beta G(M)^{\frac{2}{3}} (\rho_c)^{\frac{4}{3}} \quad ^1$$

¹See lecture notes by Onno Pols for details
https://www.astro.ru.nl/~onnop/education/stev_utrecht_notes/

Shell burning

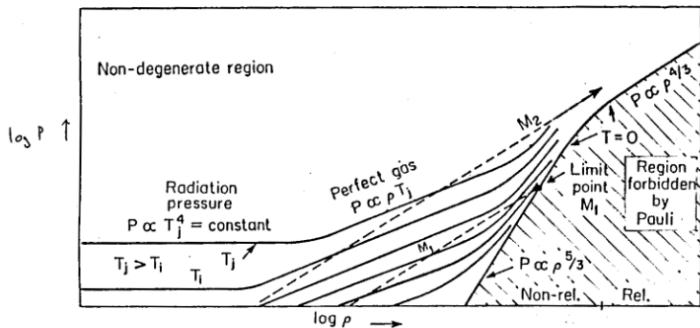
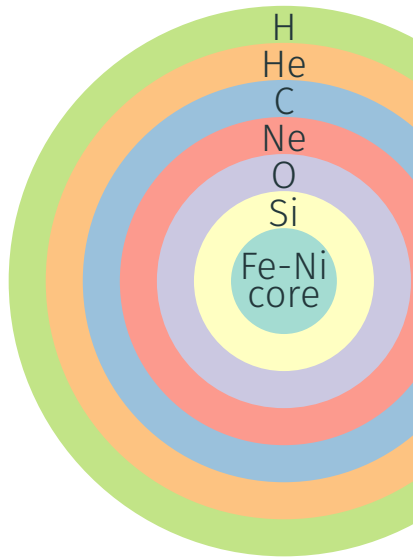


Figure 8.1. Schematic evolution in the $\log \rho$ - $\log P$ plane. Solid lines are isotherms in the equation of state; the dashed lines indicate two evolution tracks of different mass, which have a slope of $\frac{4}{3}$. See the text for an explanation.

Shell burning

- Stars of initial masses above $\sim 8M_{\odot}$
- Shell burning
- Degenerate iron core



Shell burning

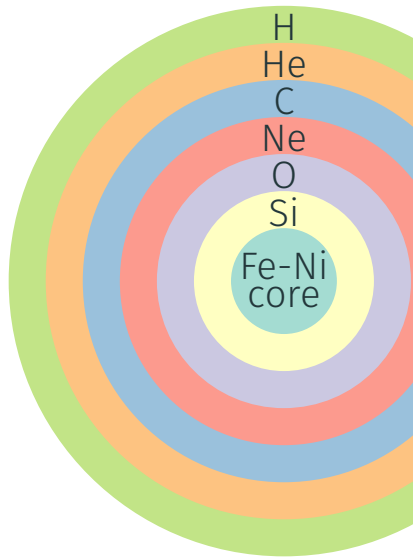
- Short burning periods
- Iron

Table 12.1. Properties of nuclear burning stages in a $15 M_{\odot}$ star (from Woosley et al. 2002).

burning stage	T (10^9 K)	ρ (g/cm^3)	fuel	main products	timescale
hydrogen	0.035	5.8	H	He	1.1×10^7 yr
helium	0.18	1.4×10^3	He	C, O	2.0×10^6 yr
carbon	0.83	2.4×10^5	C	O, Ne	2.0×10^3 yr
neon	1.6	7.2×10^6	Ne	O, Mg	0.7 yr
oxygen	1.9	6.7×10^6	O, Mg	Si, S	2.6 yr
silicon	3.3	4.3×10^7	Si, S	Fe, Ni	18 d

Shell burning: problems with this picture?

- What about multi-dimensional effects?
- Stellar evolution as a whole
- Winds, magnetic fields, and mass loss



Collapse of the Iron-core

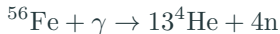
- Chandrasekhar mass limit $M_{\text{Ch}} = 5.8 Y_e^2 M_{\odot}$ [1]
- Cold ^{56}Fe white-dwarf, $M_{\text{Ch}} = 1.10618$ [3]

$$M_{\text{Ch}} = 1.44 \left(\frac{Y_e}{0.5} \right) \left[1 + \frac{s_e}{\pi Y_e} \right] [4]$$

Collapse

Collapse of the Iron-core

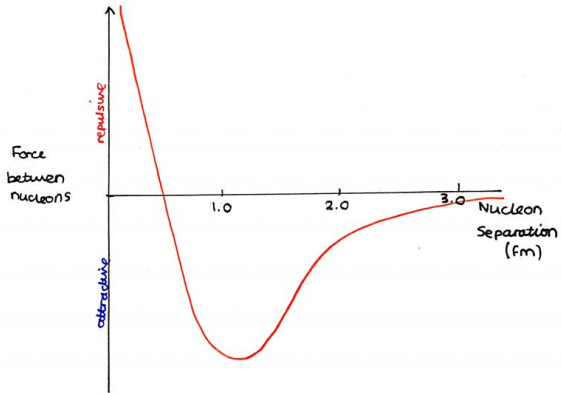
A runaway process



Continues until $\rho \sim 2 \times 10^{14} \text{ g/cm}^3$, at which point the strong nuclear force kicks in.

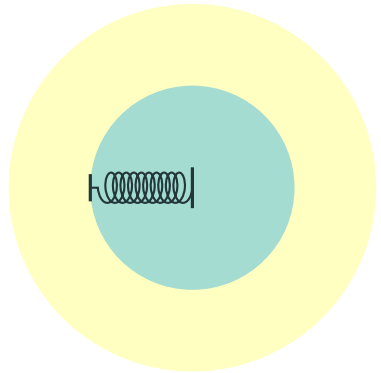
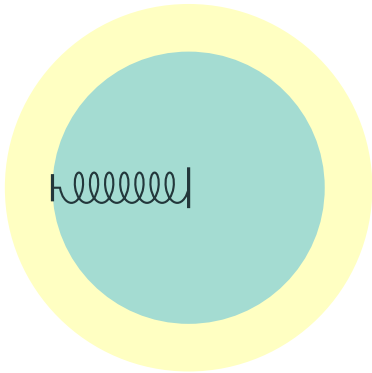
Core Bounce

- Nuclear densities
- Inner core stops contracting



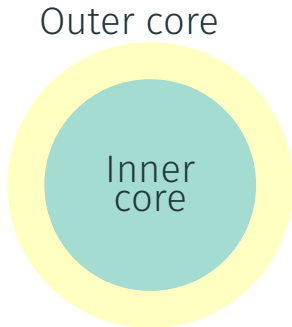
Core Recoil

The strong force acts like a spring



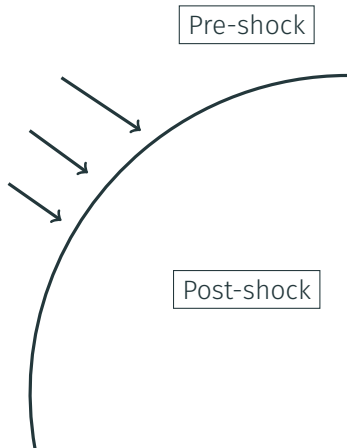
What does a compressed spring do?

- Inner core stops contracting
- Supersonically infalling outer core
- Shock wave



Shock Propagation

- Disintegration of heavy nuclei
- Density decreases
- Neutrino burst



Shock Revival

Neutrino Heating

Heating:



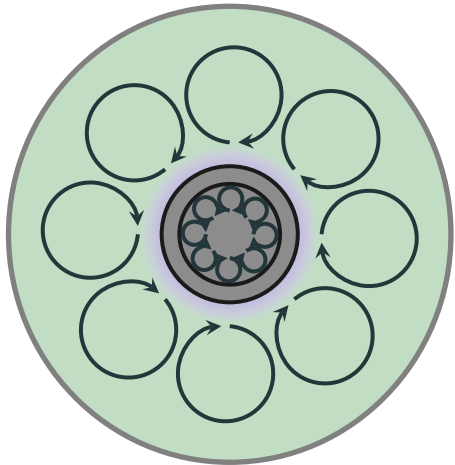
Rate $\sim 1/r^{-2}$

Cooling:



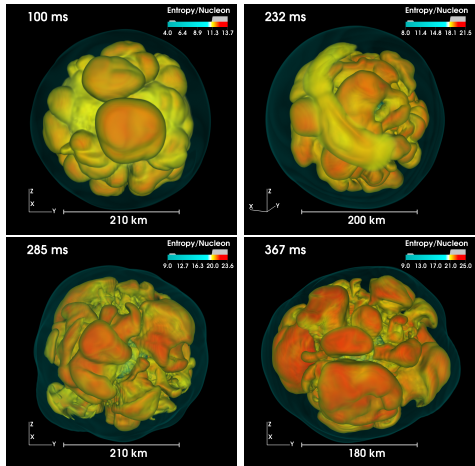
Rate $\sim 1/r^{-6}$

See [5] for a
detailed derivation



Movie time

- The standing accretion shock instability (SASI)
- Hot bubble convection
- Convection in the interior of the Proto-neutron star (PNS)



- Turbulence is key
- Material spends more time in the heating region
- One, two, and three dimensions

Energy cascade

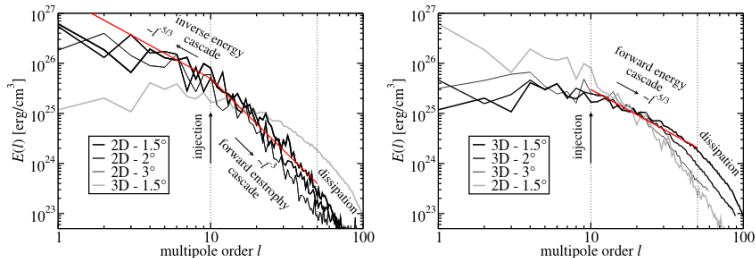


Figure 16. Turbulent energy spectra $E(l)$ as functions of the multipole order l for different angular resolution. The spectra are based on a decomposition of the azimuthal velocity v_θ into spherical harmonics at radius $r = 150$ km and 400 ms post-bounce time for $15 M_\odot$ runs with an electron-neutrino luminosity of $L_{\nu_e} = 2.2 \times 10^{52}$ erg s $^{-1}$. Left: 2D models with different angular resolution (black, different thickness) and, for comparison, the 3D model with the highest employed angular resolution (gray). Right: 3D models with different angular resolution and, for comparison, the 2D model with the highest employed angular resolution (gray). The power-law dependence and direction of the energy and enstrophy cascades (see the text) are indicated by red lines and labels for 2D models in the left panel and 3D models in the right panel. The left vertical, dotted line roughly marks the energy-injection scale, and the right vertical, dotted line denotes the onset of dissipation at high l for the best-displayed resolution.

- Proto-neutron star cools to a neutron star
- Fallback and black holes
- Nucleosynthesis and matter ejection

Numerical Simulations

- Stellar evolution
- Hydrodynamics
- General relativity
- Neutrino transport
- Equation of state
- Neutrino interactions
- Turbulence

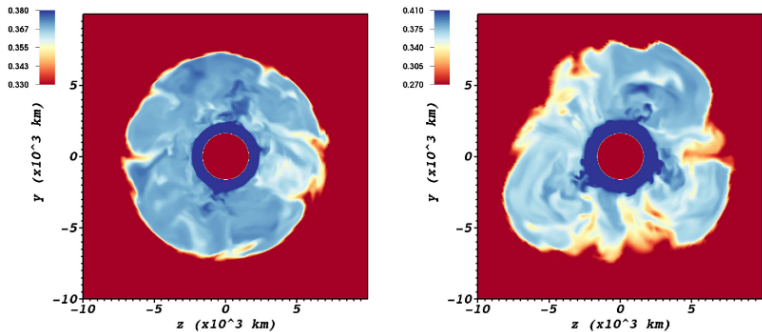


Figure 2. Slices showing the mass fraction X_{Si} of silicon at the onset of collapse in models s18-3Dr (left) and s18-3D (right). Both models are characterized by 2–3 silicon rich plumes (darker shades of blue). Due to the higher convective velocities, the boundary between the oxygen shell and the carbon shell is more strongly distorted by interfacial wave breaking in model s18-3D.)

Müller et. al. 2017 [9]

- Stellar structure
- Rotation
- Magnetic fields
- Mass loss

- Multi-dimensional
- Numerical setup and grid structure
- Turbulence

- Can you ever reach the scales you need?
- Melson et. al. 2020 [8]
- Convergence around an angular resolution of 1 degree

General Relativity

- Adds complexity and cost
- Often a “post-newtonian” approach is used [7]

$$m_{\text{TOV}} = 4\pi \int_0^r dr' r'^2 \left[\rho + e + E + \frac{vF}{r} \right]$$

$$\Phi_{\text{TOV}}(r) = 4\pi \int_r^\infty \frac{dr'}{r'^2} \left[\frac{m_{\text{TOV}}}{4\pi} + r'^3 (P + P_\nu) \right] \times \frac{1}{\Gamma} \frac{\rho + e + P}{\rho}$$

$$\Phi = \Phi - \bar{\Phi} + \Phi_{\text{TOV}}$$

- High density and temperature
- Low density regime
- Exotic particles?

Boltzmann equation

$$p^\alpha \left[\frac{\partial f_\nu}{\partial x^\alpha} - \Gamma_{\alpha\gamma}^\beta p^\gamma \frac{\partial f_\nu}{\partial p^\beta} \right] = \left[\frac{df_\nu}{d\tau} \right]_{coll}$$

Several approximations exists

- Ray-by-Ray
- Grey transport
- M1 and M2

Neutrino Interactions and Transport

Table 1 Most important neutrino processes in supernova and proto-neutron star matter.

Process	Reaction ^a
Beta-processes (direct URCA processes)	
electron and ν_e absorption by nuclei	$e^- + (A, Z) \longleftrightarrow (A, Z - 1) + \nu_e$
electron and ν_e captures by nucleons	$e^- + p \longleftrightarrow n + \nu_e$
positron and $\bar{\nu}_e$ captures by nucleons	$e^+ + n \longleftrightarrow p + \bar{\nu}_e$
“Thermal” pair production and annihilation processes	
Nucleon-nucleon bremsstrahlung	$N + N \longleftrightarrow N + N + \nu + \bar{\nu}$
Electron-positron pair process	$e^- + e^+ \longleftrightarrow \nu + \bar{\nu}$
Plasmon pair-neutrino process	$\tilde{\gamma} \longleftrightarrow \nu + \bar{\nu}$
Reactions between neutrinos	
Neutrino-pair annihilation	$\nu_e + \bar{\nu}_e \longleftrightarrow \nu_x + \bar{\nu}_x$
Neutrino scattering	$\nu_x + \{\nu_e, \bar{\nu}_e\} \longleftrightarrow \nu_x + \{\nu_e, \bar{\nu}_e\}$
Scattering processes with medium particles	
Neutrino scattering with nuclei	$\nu + (A, Z) \longleftrightarrow \nu + (A, Z)$
Neutrino scattering with nucleons	$\nu + N \longleftrightarrow \nu + N$
Neutrino scattering with electrons and positrons	$\nu + e^\pm \longleftrightarrow \nu + e^\pm$

^a N means nucleons, i.e., either n or p , $\nu \in \{\nu_e, \bar{\nu}_e, \nu_\mu, \bar{\nu}_\mu, \nu_\tau, \bar{\nu}_\tau\}$, $\nu_x \in \{\nu_\mu, \bar{\nu}_\mu, \nu_\tau, \bar{\nu}_\tau\}$

See [6]

Neutrino Interactions and Transport

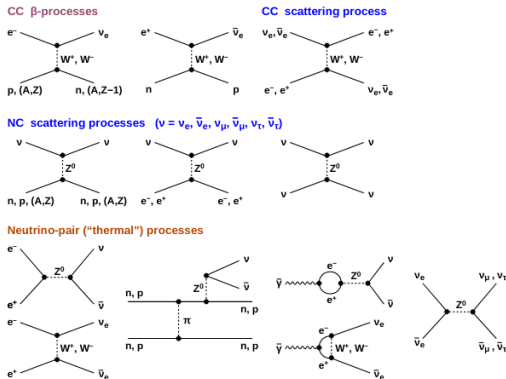


Fig. 3 Feynman diagrams for the lowest-order contributions to the most relevant neutrino interactions in supernova cores. Charged-current (CC) reactions are mediated by W^\pm bosons, neutral-current (NC) reactions by electrically neutral Z^0 bosons. The charged-current β -processes are responsible for the production and absorption of ν_e and $\bar{\nu}_e$ by lepton-capture reactions on nucleons (*top row, left*). Scattering processes include the charged-current interactions of ν_e and $\bar{\nu}_e$ with electrons and positrons (*top row, right*) and neutral-current scatterings of neutrinos and antineutrinos of all flavors with nuclei, neutrons, protons, electrons, positrons, and neutrinos (*middle row*). Neutrino-pair processes are responsible for the creation and annihilation of neutrino-antineutrino pairs of all flavors. They include electron-positron pair annihilation through neutral and charged currents, nucleon bremsstrahlung, the charged- and neutral-current plasmon-neutrino processes, and neutrino-pair conversion between different flavors (*bottom row, from left to right*).

TABLE I. Neutrino reactions with muons.

$\nu + \mu^- \rightleftharpoons \nu' + \mu'^-$	$\nu + \mu^+ \rightleftharpoons \nu' + \mu'^+$
$\nu_\mu + e^- \rightleftharpoons \nu_e + \mu^-$	$\bar{\nu}_\mu + e^+ \rightleftharpoons \bar{\nu}_e + \mu^+$
$\nu_\mu + \bar{\nu}_e + e^- \rightleftharpoons \mu^-$	$\bar{\nu}_\mu + \nu_e + e^+ \rightleftharpoons \mu^+$
$\bar{\nu}_e + e^- \rightleftharpoons \bar{\nu}_\mu + \mu^-$	$\nu_e + e^+ \rightleftharpoons \nu_\mu + \mu^+$
$\nu_\mu + n \rightleftharpoons p + \mu^-$	$\bar{\nu}_\mu + p \rightleftharpoons n + \mu^+$

See [2]

- Many different things goes into the simulations
- Uncertainties in almost every aspect
- Approximations are necessary

Performing Simulations

- GR1D: <https://github.com/evanoconnor/GR1D>
- NuLiB: <https://github.com/evanoconnor/NuLib>
- FLASH: <http://flash.uchicago.edu/site/flashcode/>
- Einstein Toolkit: <https://einsteintoolkit.org>
- Specter:
<https://github.com/sxs-collaboration/spectre>
- Quokka: <https://github.com/BenWibking/quokka>

- GR1D
- A $75 M_{\odot}$ star
- Polytropic equation of state

```
git clone https://github.com/evanoconnor/GR1D.git  
cd GR1D
```


Setting up GR1D

F90=

F90FLAGS=-g -O3 -fopenmp

LDLFLAGS=-g -O3 -fopenmp

MODINC="-I./"

HDF5DIR=

HDF5INCS=-I\${HDF5DIR}/include

HDF5LIBS=-L\${HDF5DIR}/lib -lhdf5 -lhdf5_fortran -lhdf5 -lz

LAPACKDIR=

LAPACKLIBS=-L\${LAPACKDIR}/lib -llapack -lblas

HAVE_NUC_EOS=1

HAVE_RESTART=1

HAVE_LEAK_ROS=0

HAVE_LAPACK=1

Setting up GR1D

```
cp sample_parameter_files/latest_recommended_params ./parameters
```

```
##### Job parameters #####  
jobname = "BHF"  
GR = 1 # 1 for GR, 0 for Newtonian  
outdir = "u70-Black-hole"  
initial_data = "Collapse"  
profile_name = "./u70.short" # stellar profile  
WHW02profile = 1  
profile_type = 1  
gravity_active = 1 # do we want gravity?  
ntmax = 100000000 # maximum timestep  
tend = 0.7d0 # maximum time
```

Setting up GR1D

```
##### Grid parameters #####  
geometry = 2 # 1: planar, 2: spherical  
gridtype = "custom" # "log", "unigrid", "custom", "custom2"  
### custom & custom2 input parameters ###  
grid_custom_rad1 = 20.0d5  
grid_custom_dx1 = 3.0d4 # smallest radial zone  
  
rmax_from_profile = 1 # take rmax to be where rho=3.0d3g/cm^3  
rho_cut = 2.0d3 # density to cut profile at if rmax_from_profile = 1  
  
radial_zones = 600 # number of radial zones  
ghosts1 = 4 # number of ghost cells
```

Setting up GR1D

```
##### EOS parameters #####  
eoskey = 1 # hybrid: 1 # poly: 2 # hot nuclear: 3 # ideal: 4  
eos_table_name = ""  
hybridgamma_th = 1.30d0 # hybrid gamma_th  
hybridgamma1 = 1.31d0 # hybrid gamma_th  
hybridgamma2 = 2.40d0 # hybrid gamma_th
```

Setting up GR1D

```
##### M1 settings #####
```

```
do_M1 = 0 #1 for M1 transport scheme
```

```
##### Neutrino parameters #####
```

```
fake_neutrinos = 0 # 1 for ANY fake neutrino scheme
```

Running GR1D

Finding envelope binding energy

Binding energy of envelope (ergs): 3.4481747535627853E+051

Finding accretion radii

Setting up PPM coefficients

Done with initial data :-)

Begin time integration loop:

0	0.000000E+00	1.050000E-07	1.368096E+09	9.934271E-01
100	2.740526E-04	1.380763E-05	1.385673E+09	9.934264E-01
200	2.474790E-03	2.256444E-05	1.383073E+09	9.934198E-01
300	4.731547E-03	2.253527E-05	1.388112E+09	9.934114E-01
400	6.980886E-03	2.244709E-05	1.401145E+09	9.934012E-01
500	9.220404E-03	2.233893E-05	1.422465E+09	9.933892E-01

```
cd u75-bh  
ls
```

```
M_innercore.dat      eps_kin.xg           press.xg  
W.xg                 list_of_restartfiles.txt  pressth.xg  
X.xg                 mass_bary.xg         r_Xmax.dat  
accreted_mass.dat   mass_grav.xg         r_rho1e11.dat  
accretion_rates.dat m_bary_Xmax.dat      r_rho1e12.dat  
alpha.xg             m_bary_rho1e12.dat   ye.xg  
alpha_c_t.dat        m_bary_shock.dat     ye_c_t.dat  
cs.xg                mgrav_Xmax.dat        ynu_c_t.dat  
c_sound_c_t.dat     mgrav_rho1e12.dat    rho.xg  
dyedt_hydro_c_t.dat mgrav_shock.dat       rho_c_t.dat  
eps.xg               parameters            time_c.dat
```

Different output files

- .xg : Radial profiles at several time-steps
- .dat : Special data versus time
- _c.dat Values at the central zone versus time

The structure of the .xg files depend on the grid and the output frequency can be changed in the parameter file.

Exercise

- Repeat the evolution
- Tabulated EOS
- M1 transport
- Bounce time, black hole formation, difference in the evolution?

References

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- [3] Kuantay Boshkayev, Jorge A. Rueda, Remo Ruffini, and Ivan Siutsou. On General Relativistic Uniformly Rotating White Dwarfs. *ApJ*, 762(2):117, January 2013. doi: 10.1088/0004-637X/762/2/117.

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- [5] H. Th. Janka. Conditions for shock revival by neutrino heating in core-collapse supernovae. *A&A*, 368:527–560, March 2001. doi: 10.1051/0004-6361:20010012.
- [6] Hans-Thomas Janka. *Neutrino Emission from Supernovae*, page 1575. 2017. doi: 10.1007/978-3-319-21846-5_4.

- [7] Marek, A., Dimmelmeier, H., Janka, H.-Th., Müller, E., and Buras, R. Exploring the relativistic regime with newtonian hydrodynamics: an improved effective gravitational potential for supernova simulations. *A&A*, 445(1):273–289, 2006. doi: 10.1051/0004-6361:20052840. URL <https://doi.org/10.1051/0004-6361:20052840>.
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- [9] Bernhard Müller, Tobias Melson, Alexander Heger, and Hans-Thomas Janka. Supernova simulations from a 3D progenitor model – Impact of perturbations and evolution of explosion properties. *Monthly Notices of the Royal Astronomical Society*, 472 (1):491–513, 08 2017. ISSN 0035-8711. doi: 10.1093/mnras/stx1962. URL <https://doi.org/10.1093/mnras/stx1962>.